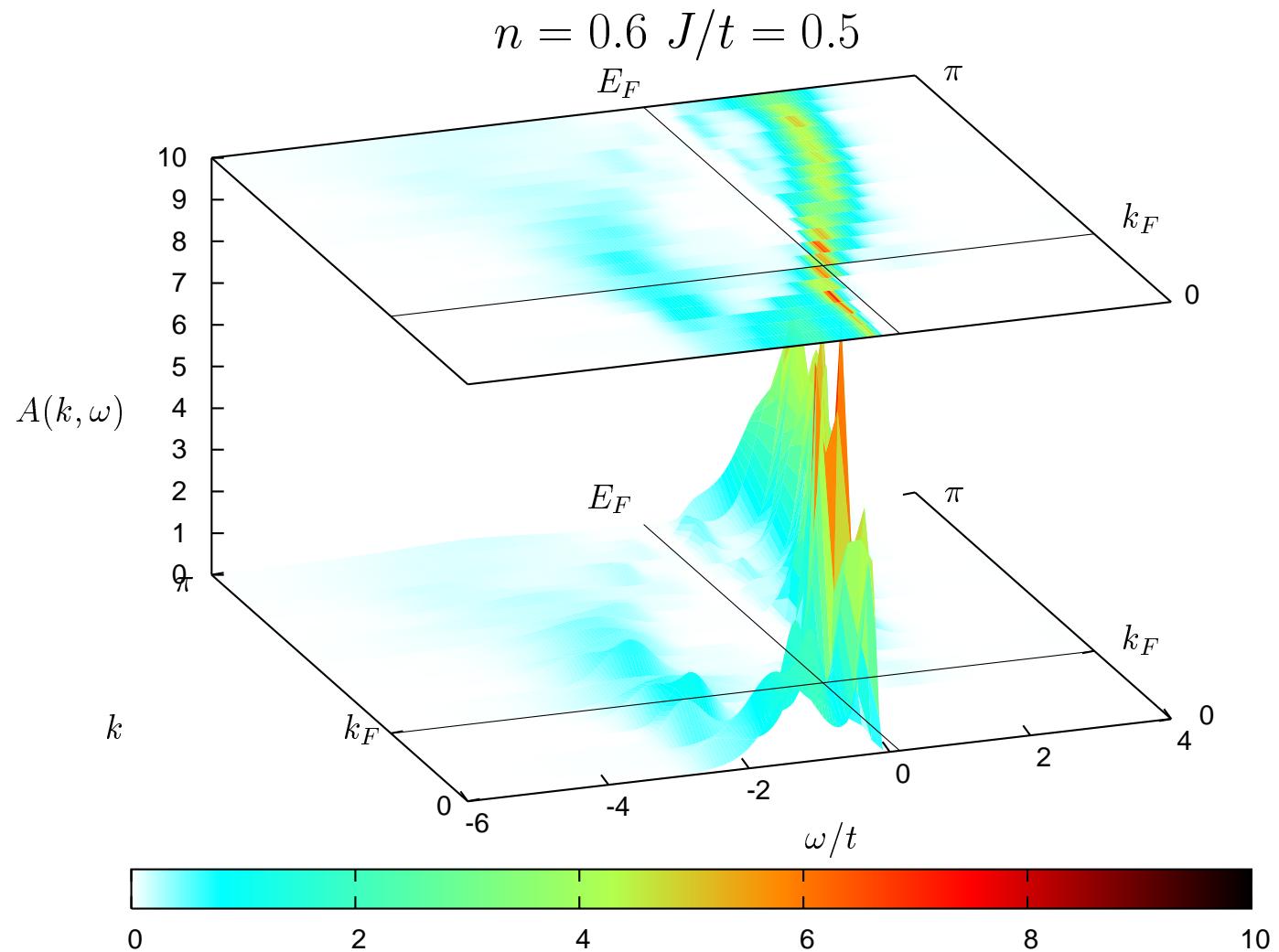


Monte Carlo simulations of quantum systems with global updates

Alejandro Muramatsu
Institut für Theoretische Physik III
Universität Stuttgart

Spinons, holons, and antiholons in one dimension

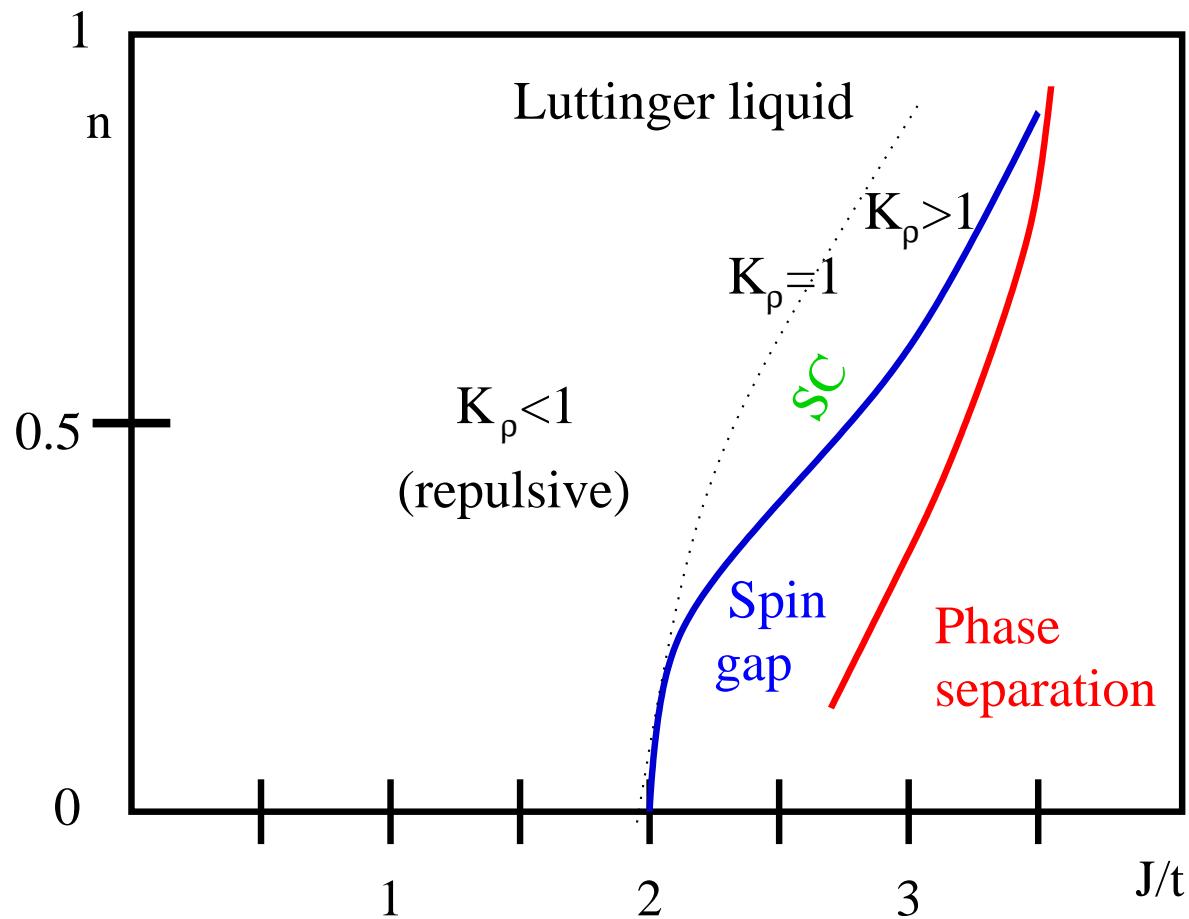


5.1 The t-J model in one dimension

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Phase diagram

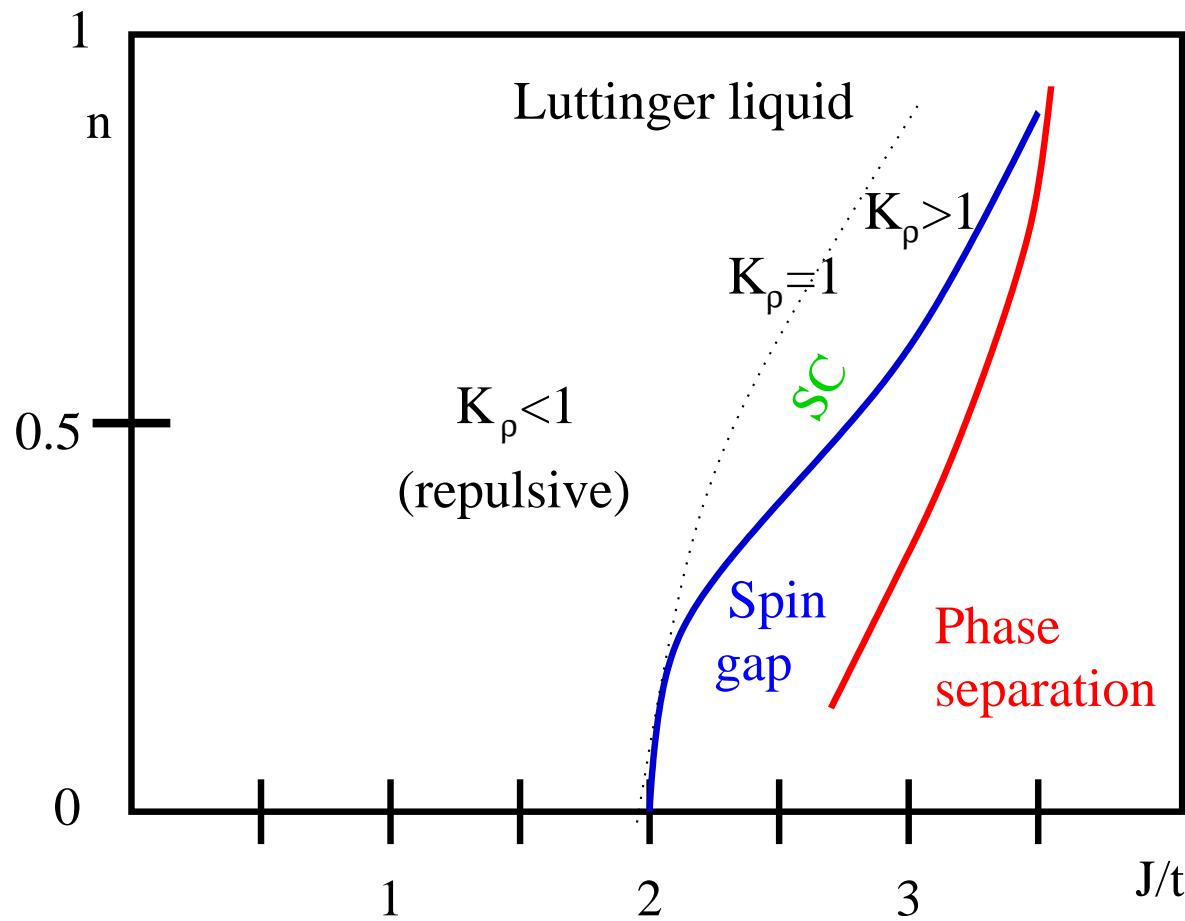
M. Ogata, M.U. Luchini, S. Sorella, and F.F. Assaad, Phys. Rev. Lett. **66**, 2388 (1991)



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Phase diagram

M. Ogata, M.U. Luchini, S. Sorella, and F.F. Assaad, Phys. Rev. Lett. **66**, 2388 (1991)



Phase diagram, with phases similar to those in high T_c superconductors

Exact results for the nearest neighbor t-J model

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■ At $J = 0$

One-particle spectral function from Ogata-Shiba wavefunction

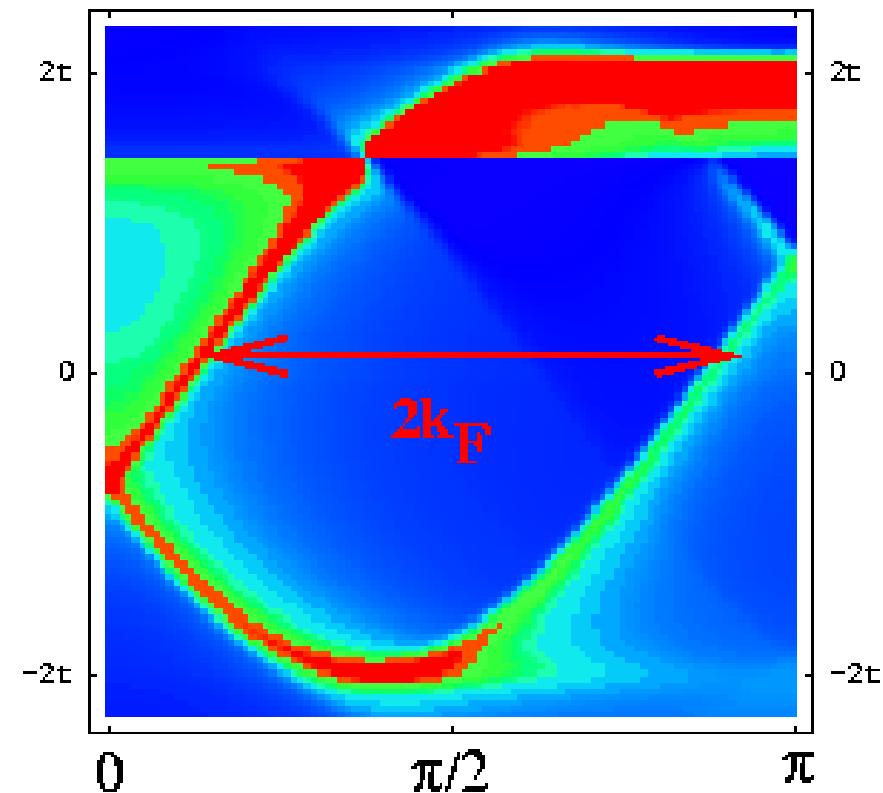
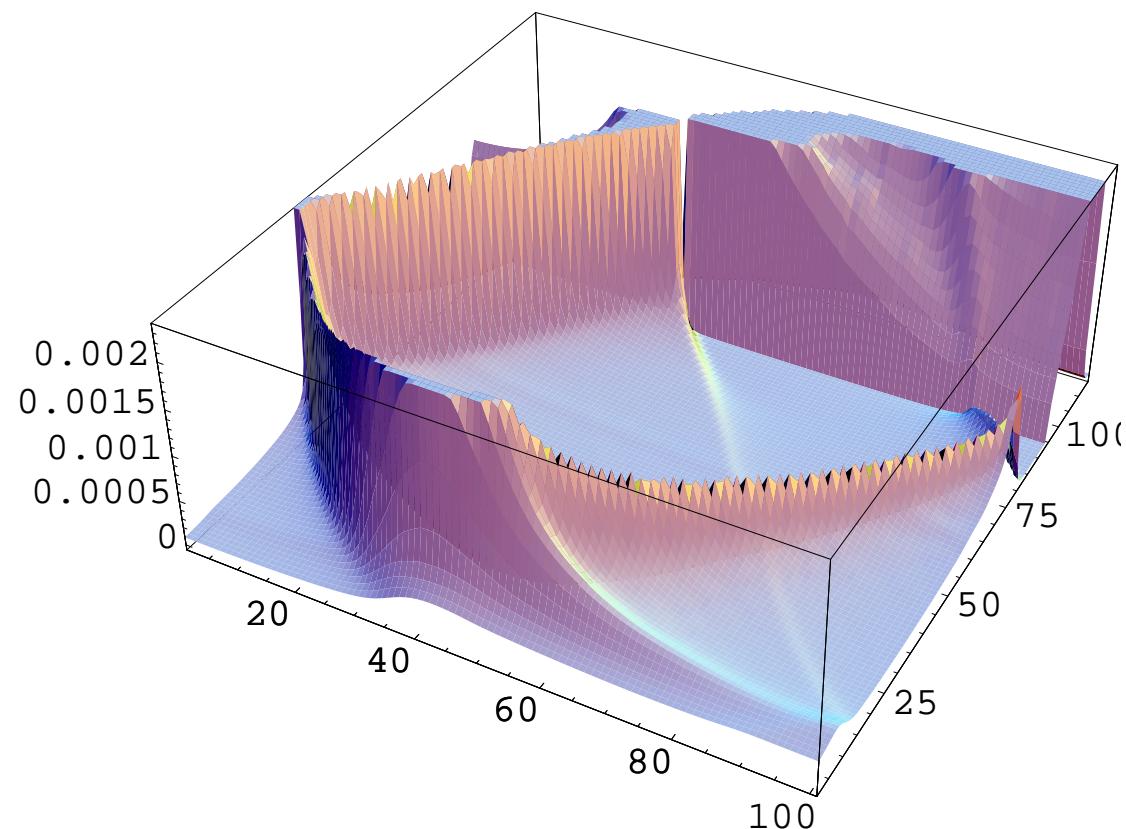
K. Penc, K. Hallberg, F. Mila, and H. Shiba, Phys. Rev. Lett. **77**, 1390 (1996)

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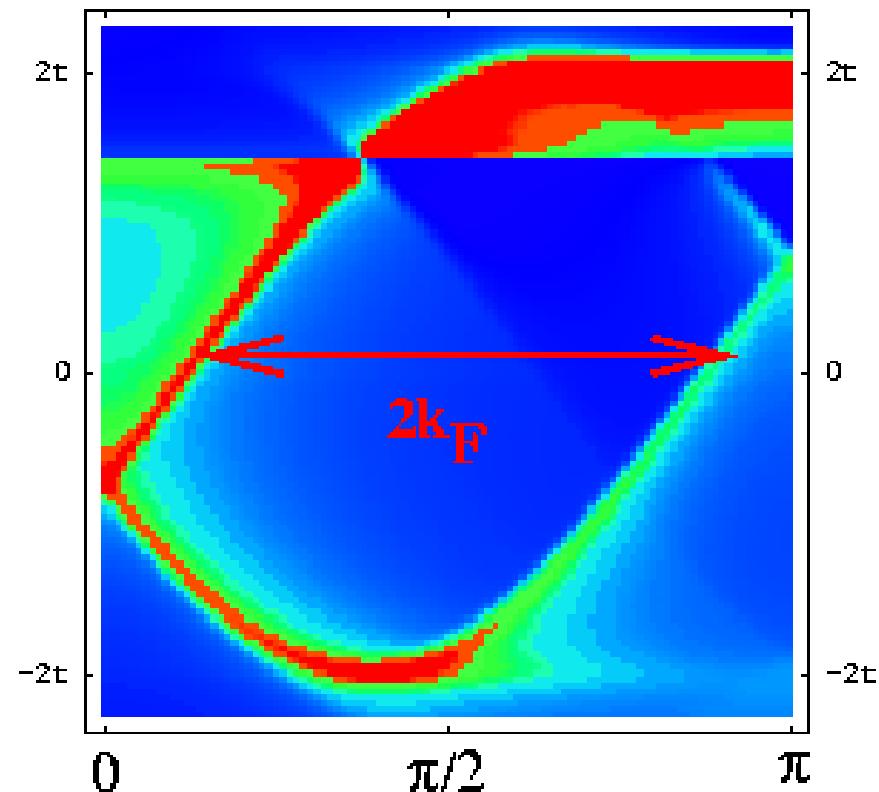
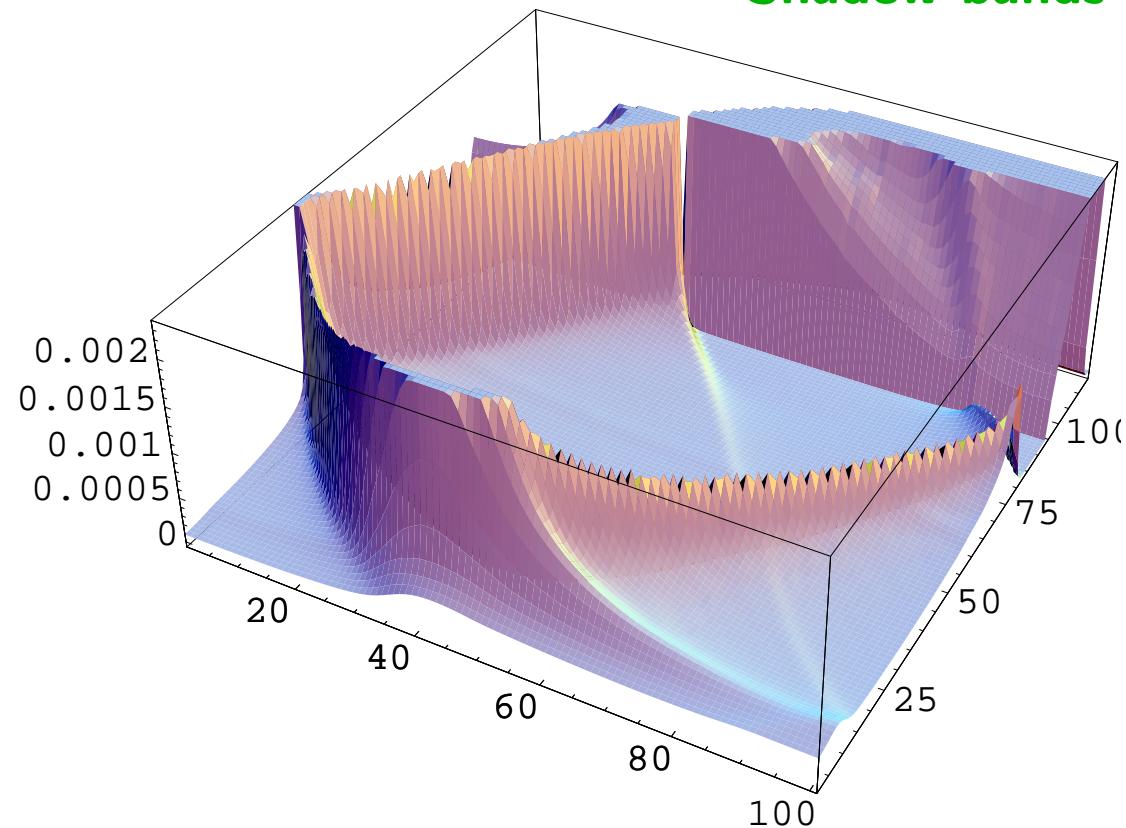
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Shadow bands

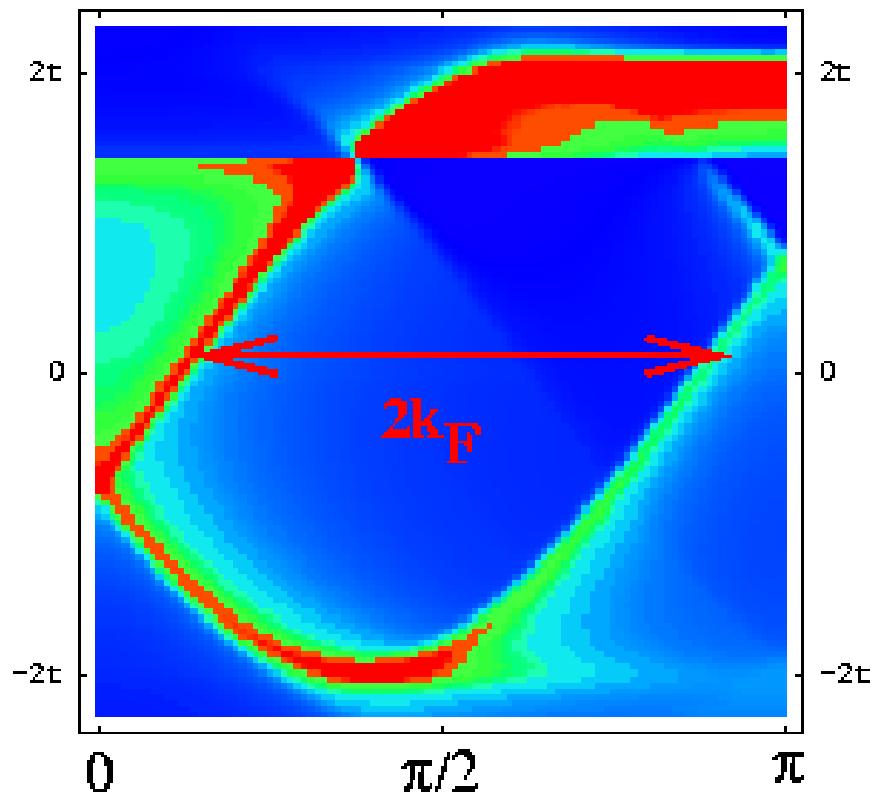
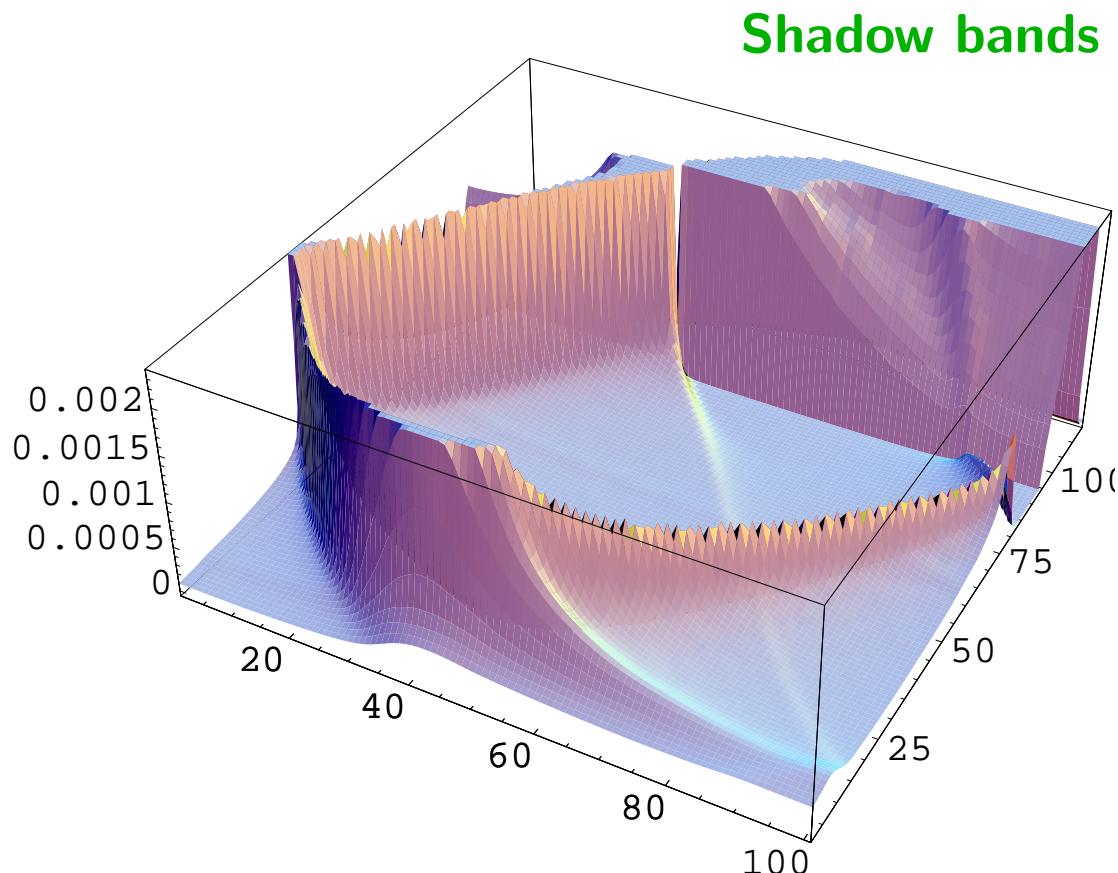


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■ At $J = 2t$

Bethe-Ansatz solution

P.A. Bares and G. Blatter, Phys. Rev. Lett. **64**, 2567 (1990)

No correlations functions

5.2 The $1/r^2$ (inverse square) t-J model in one dimension

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$$H_{IS-t-J} = - \sum_{i < j, \sigma} t_{ij} (\tilde{c}_{i,\sigma}^\dagger \tilde{c}_{j,\sigma} + \text{h.c.}) + \sum_{i < j} J_{ij} \left(\mathbf{S}_i \cdot \mathbf{S}_j - \frac{1}{4} \tilde{n}_i \tilde{n}_j \right),$$

where at the supersymmetric (SuSy) point $J = 2t$,

$$t_{ij} = J_{ij}/2 = t(\pi/L)^2 / \sin^2(\pi(i-j)/L)$$

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Compact support and excitation content

Z.N.C. Ha and F.D.M. Haldane, Phys. Rev. Lett. **73**, 2887 (1994)

spinon	$Q=0$,	$S=1/2$	semion
holon	$Q=+e$,	$S=0$	semion
antiholon	$Q=-2e$,	$S=0$	boson

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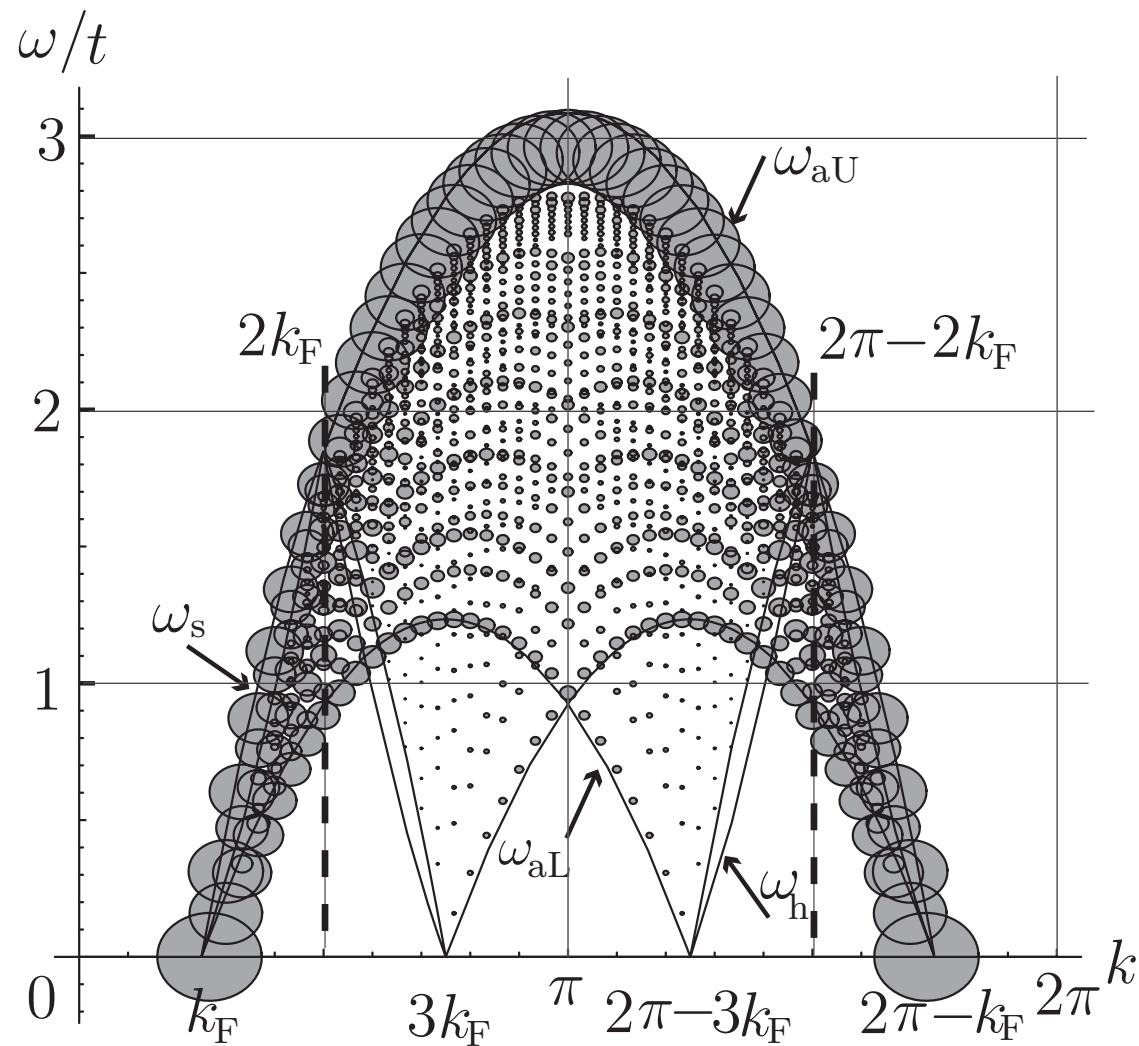
Semion $\rightarrow s_j s_\ell = i s_\ell s_j \rightarrow$ ‘half a fermion’

Spectral function for electron addition

M. Arikawa, Y. Saiga, and Y. Kuramoto, Phys. Rev. Lett. **86**, 3096 (2001)

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Analytic expressions

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Spectral function for electron addition ($0 \leq k < 2\pi$)

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$$A^+(k, \omega) = A_R(k, \omega) + A_L(k, \omega) + A_U(k, \omega)$$

where

$$\begin{aligned} A_R(k, \omega) = & \frac{1}{4\pi} \int_0^{k_F} dq_h \int_0^{k_F - q_h} dq_s \int_0^{2\pi - 4k_F} dq_a \delta(k - k_F - q_s - q_h - q_a) \\ & \times \delta[\omega - \epsilon_s(q_s) - \epsilon_h(q_h) - \epsilon_a(q_a)] \frac{\epsilon_s^{g_s-1}(q_s) \epsilon_h^{g_h-1}(q_h) \epsilon_a^{g_a-1}(q_a)}{(q_h + q_a/2)^2}, \end{aligned}$$

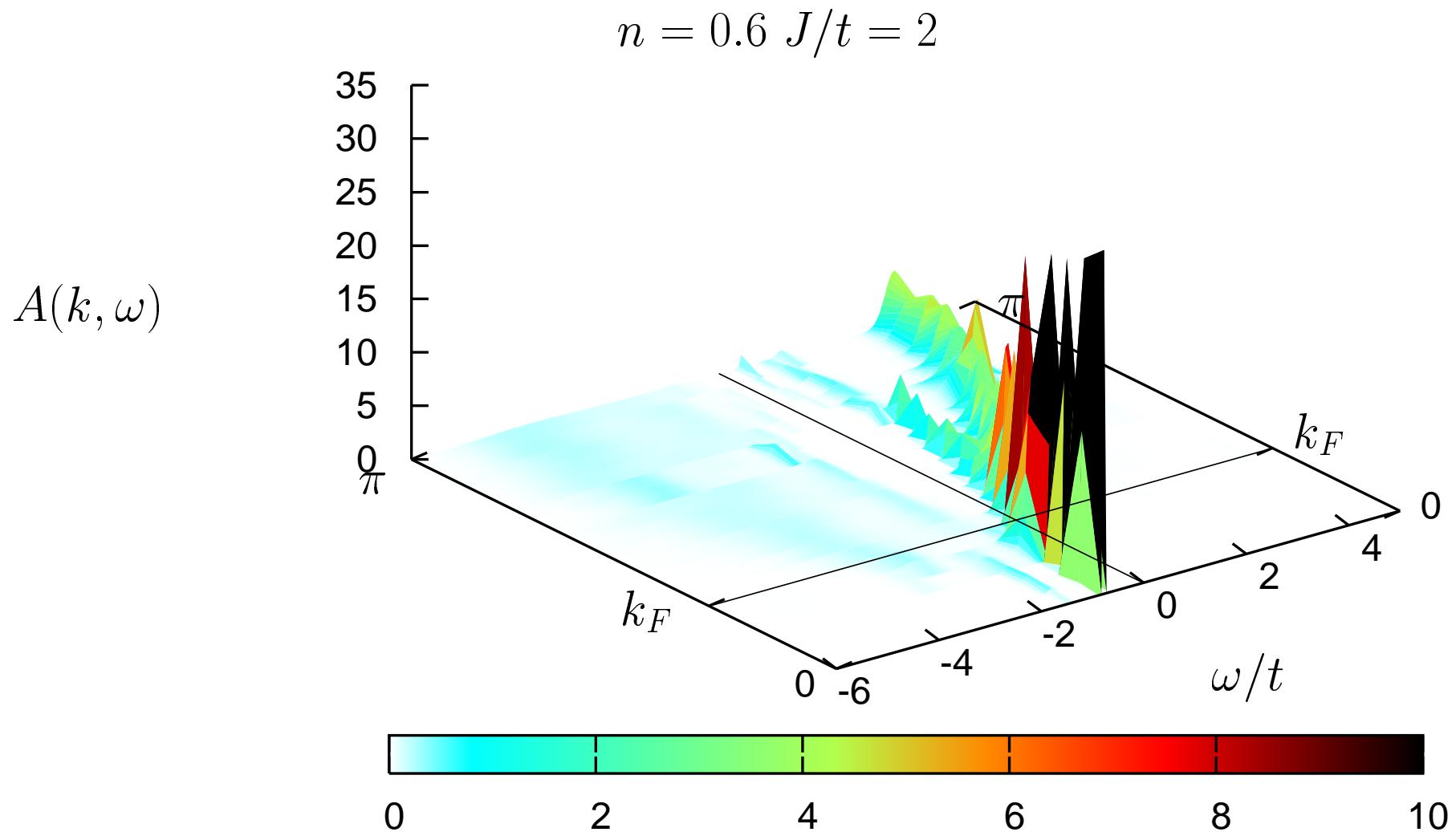
with $A_L(k, \omega) = A_R(2\pi - k, \omega)$

$g_s = 1/2$, $g_h = 1/2$, **and** $g_a = 2$, **statistical parameters, and**

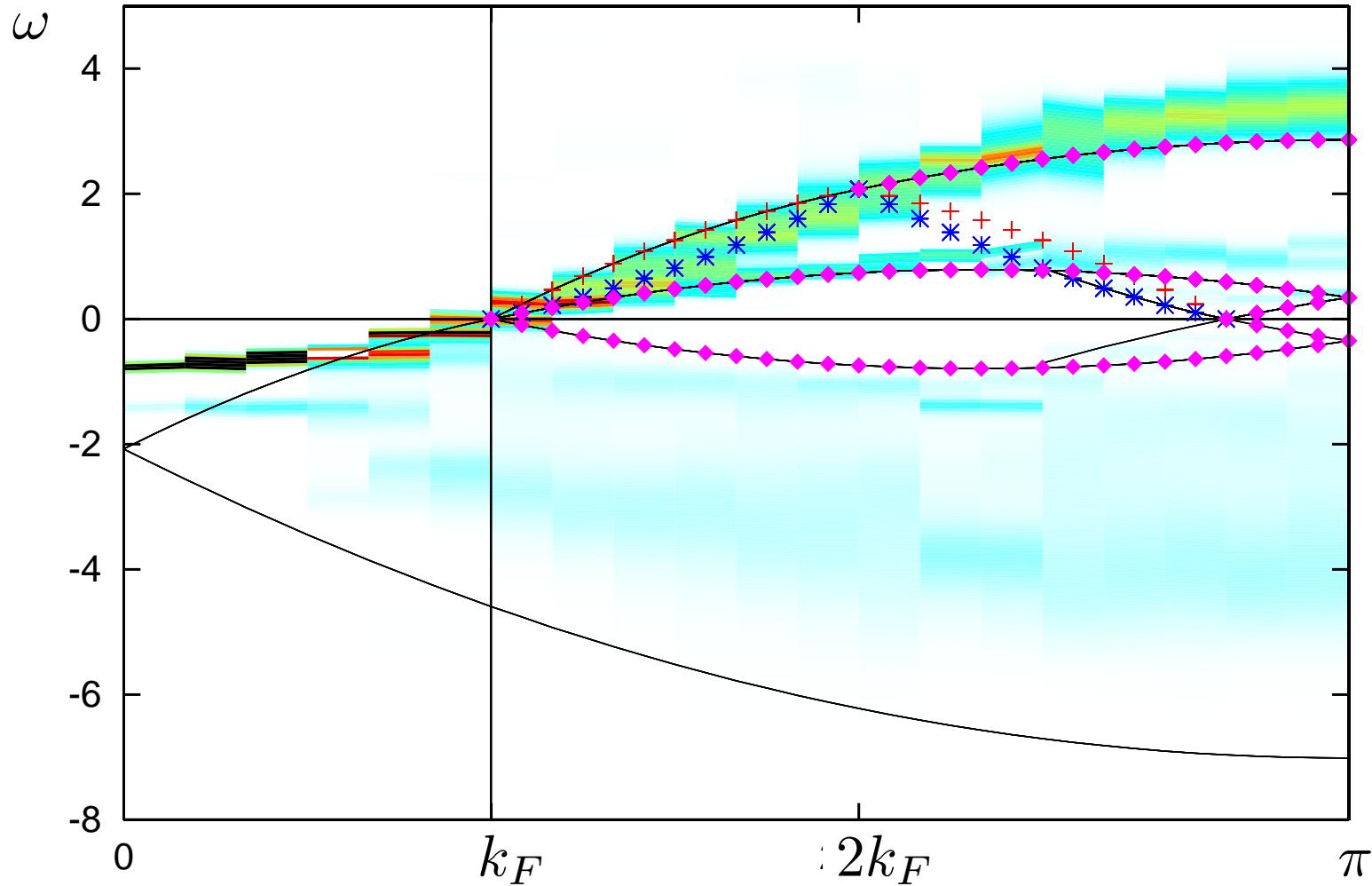
$$A_U(k, \omega) = \sqrt{\frac{\epsilon_a(k - 2k_F)}{k(\pi - k/2)}} \delta\left\{\omega - [\epsilon_s(k_F) + \epsilon_a(k - 2k_F)]\right\}, \quad (2k_F \leq k \leq 2\pi - 2k_F)$$

5.3 Spectral functions for the n.n. t-J model from QMC

C. Lavalle, M. Arikawa, S. Capponi, F.F. Assaad, and A. Muramatsu,
PRL **90**, 216401, (2003)



Spinon, holons, and antiholons in the n.n. t-J model at $J=2t$



+ spinon:	$\epsilon_{sR(L)}(q_s)$	=	$tq_s (\pm v_s^0 - q_s)$	$0 \leq q_s \leq k_F$
* holon:	$\epsilon_{hR(L)}(q_h)$	=	$tq_h (q_h \pm v_c^0)$	$0 \leq q_h \leq k_F$
♦ antiholon:	$\epsilon_{\bar{h}}(q_{\bar{h}})$	=	$tq_{\bar{h}}(2v_c^0 - q_{\bar{h}})/2$	$0 \leq q_{\bar{h}} \leq 2\pi - 4k_F$

$$v_s^0 = \pi, v_c^0 = \pi(1 - n)$$

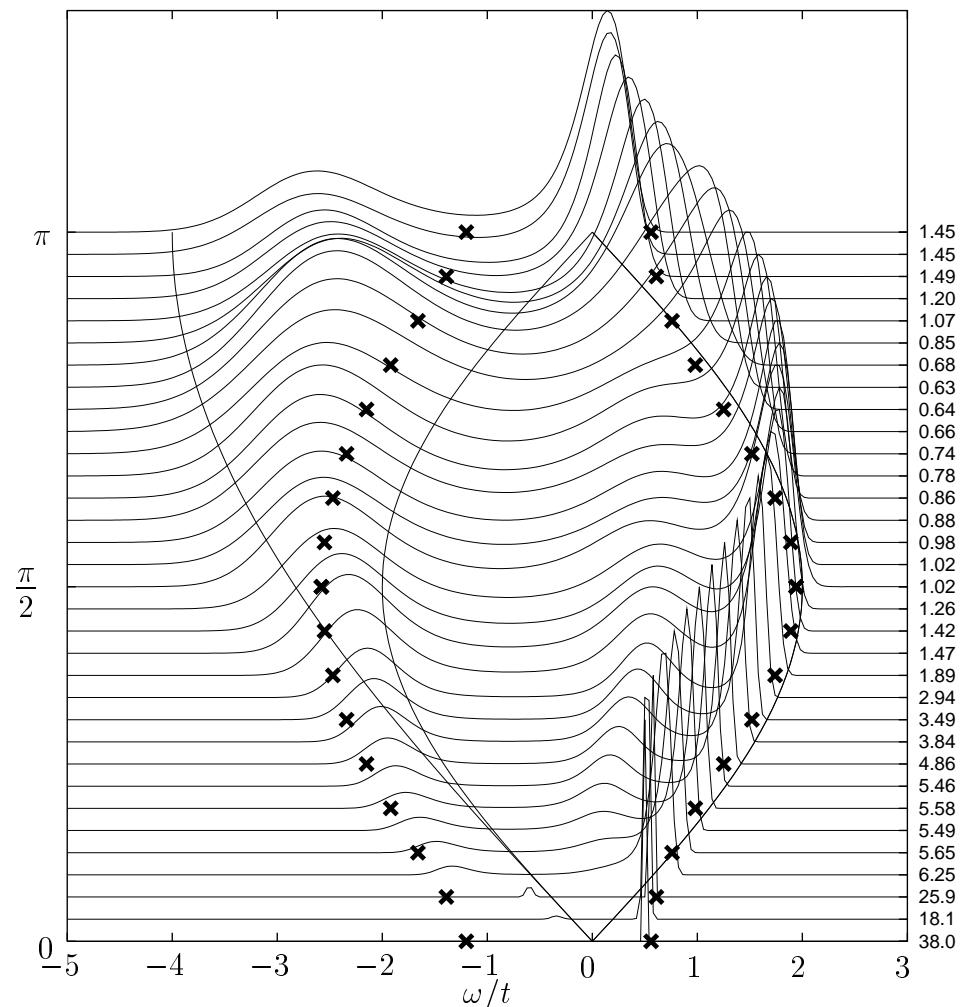
Excitation content of the hole spectrum at the supersymmetric point

M. Brunner, F.F. Assaad, and A. Muramatsu, Eur. Phys. J. B **16**, 209 (2000)

Supersymmetric point $J/t = 2 \rightarrow$ exact holon and spinon dispersions from Bethe-Ansatz

P.-A. Bares and G. Blatter, Phys. Rev. Lett. **64**, 2567 (1990)

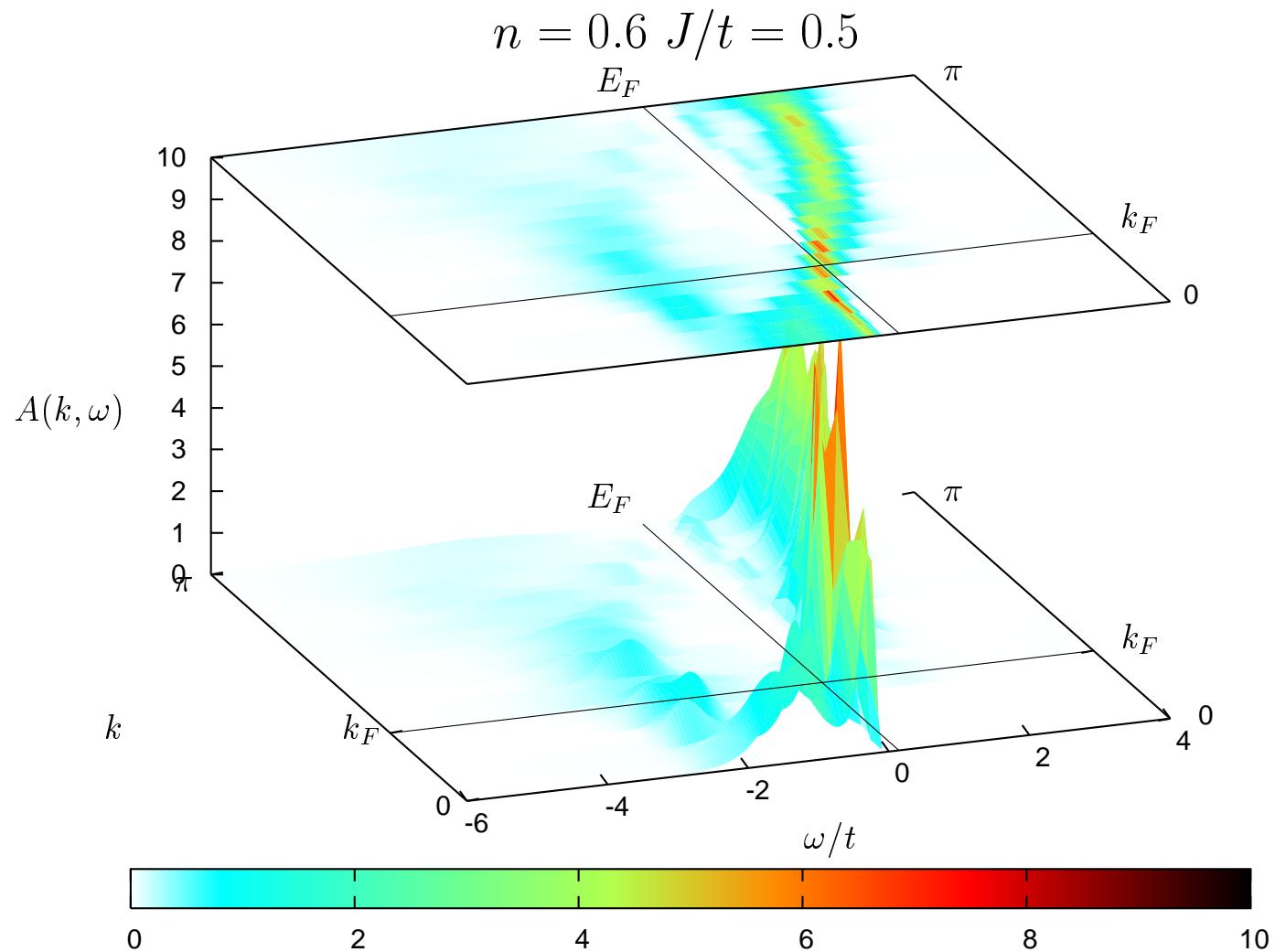
P.-A. Bares, G. Blatter, and M. Ogata, Phys. Rev. B **44**, 130 (1991)



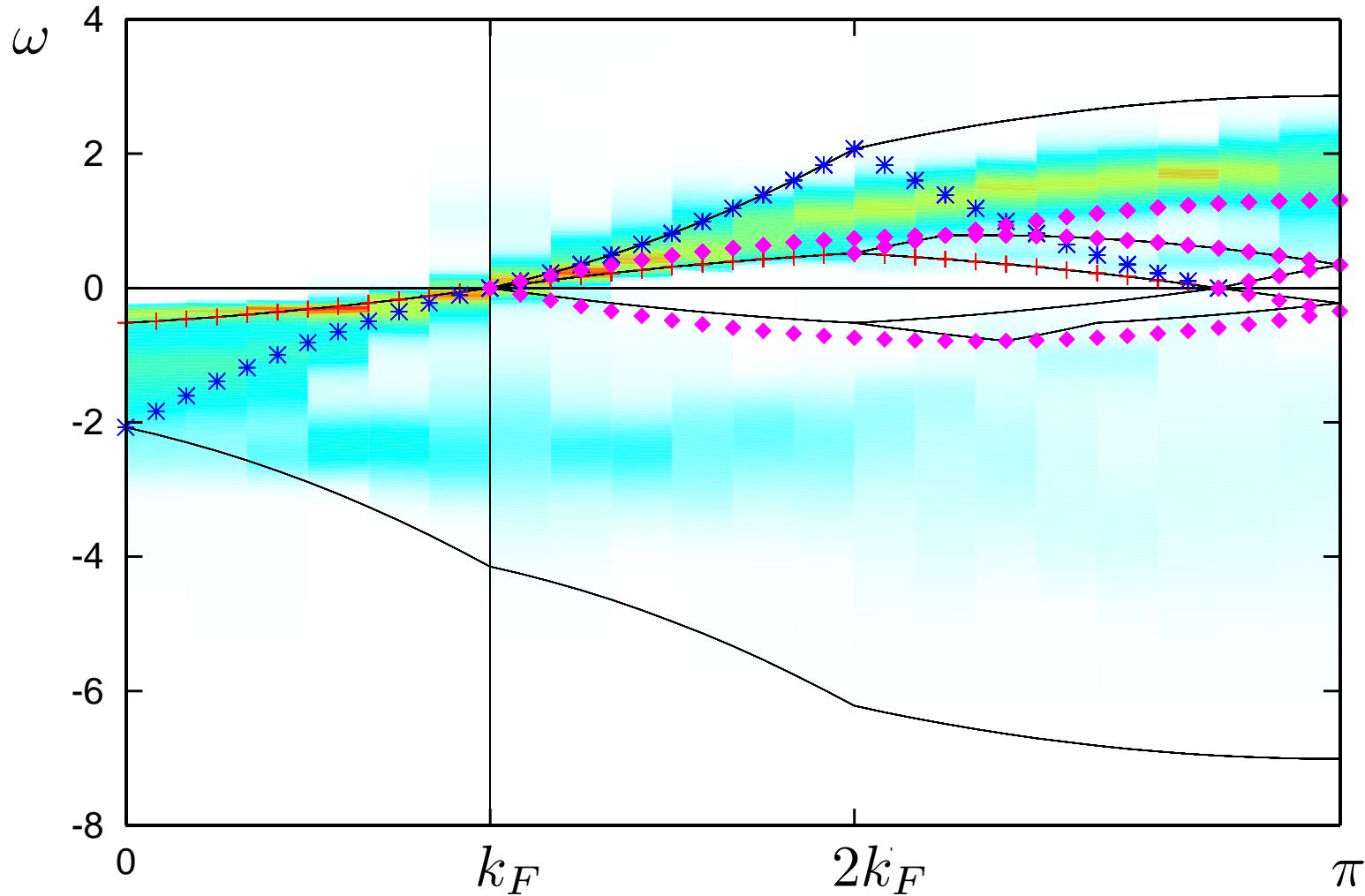
Full line: One holon + one spinon with dispersions given by charge-spin separation Ansatz.

Crosses: One holon + one spinon with dispersions given by Bethe-Ansatz

Spectral function at $J = 0.5 t$



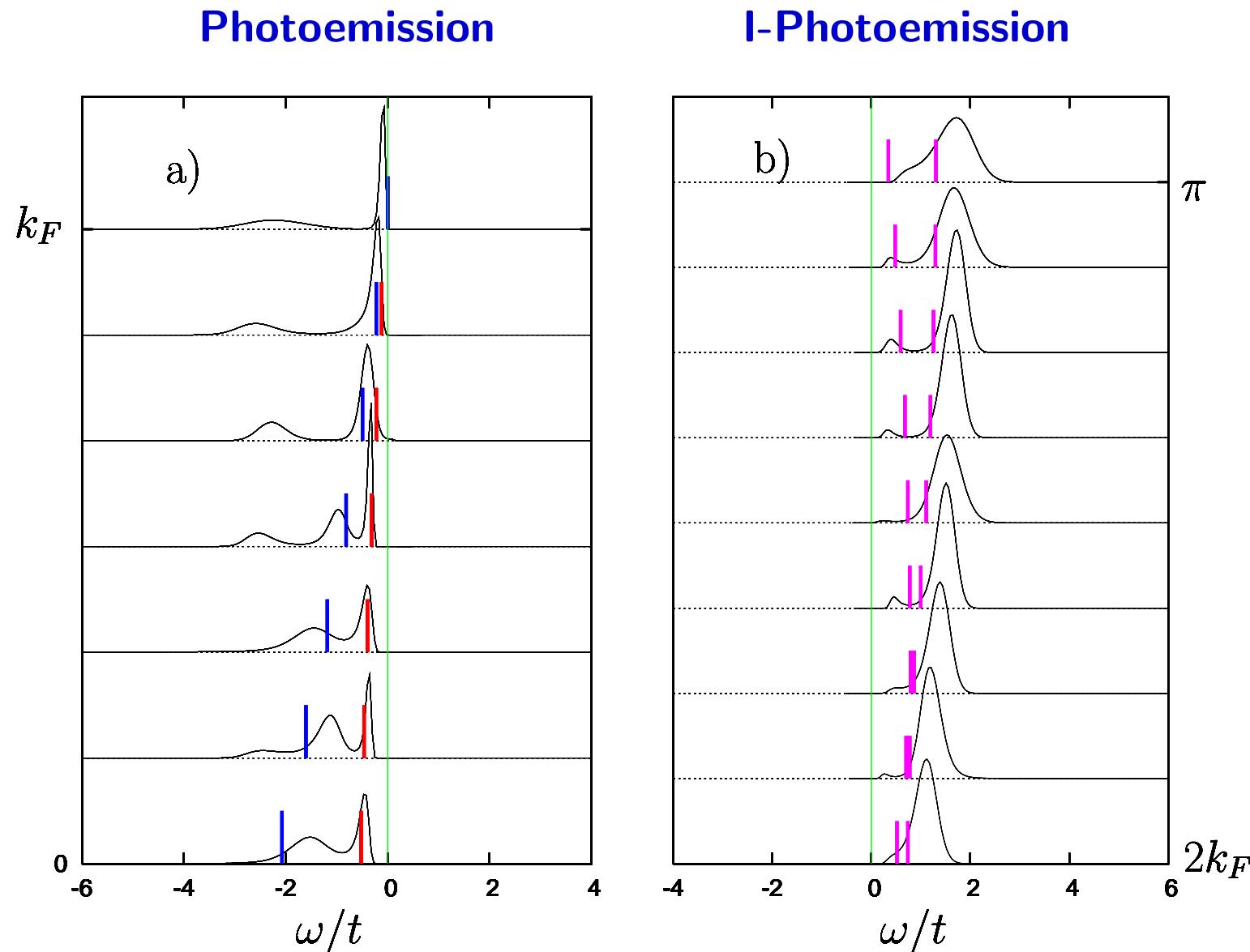
Spinon, holons, and antiholons at $J = 0.5 t$ and $n = 0.6$



+ spinon:	$\epsilon_{sR(L)}(q_s)$	=	$(J/2) q_s (\pm v_s^0 - q_s)$	$0 \leq q_s \leq k_F$
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♦ antiholon:	$\epsilon_{\bar{h}}(q_{\bar{h}})$	=	$t q_{\bar{h}} (2v_c^0 - q_{\bar{h}})/2$	$0 \leq q_{\bar{h}} \leq 2\pi - 4k_F$

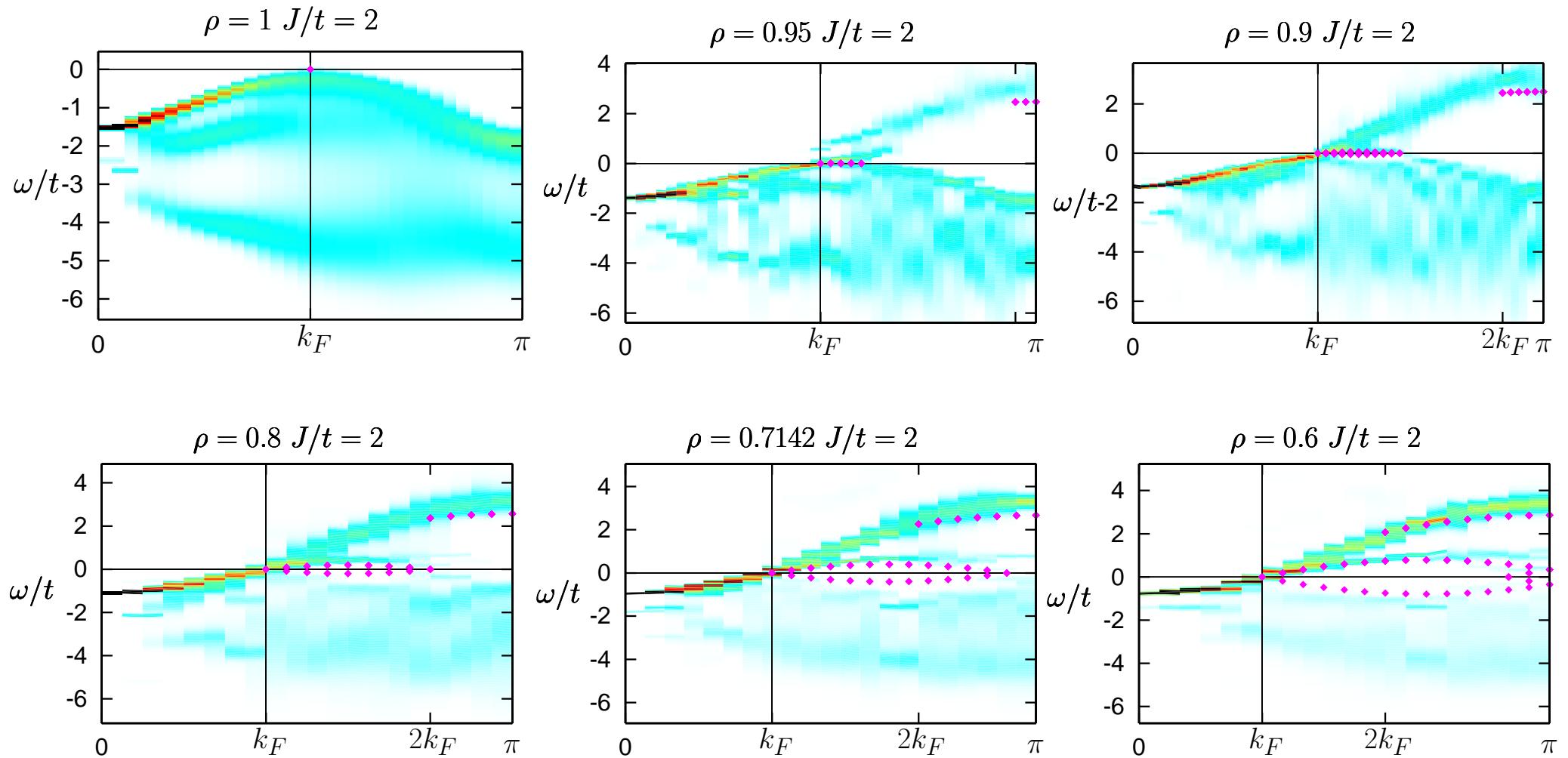
$$v_s^0 = \pi, v_c^0 = \pi(1 - n)$$

Charge-spin separation at $J = 0.5 t$



Antiholons at $J = 0.5 t$ vs. doping

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Antiholons are generic to the nearest neighbor t-J model at finite J

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- **New excitation with $Q = 2e$, $S = 0$, and $m_{ah} = 2m_h$**
→ **not charge conjugated to holons**

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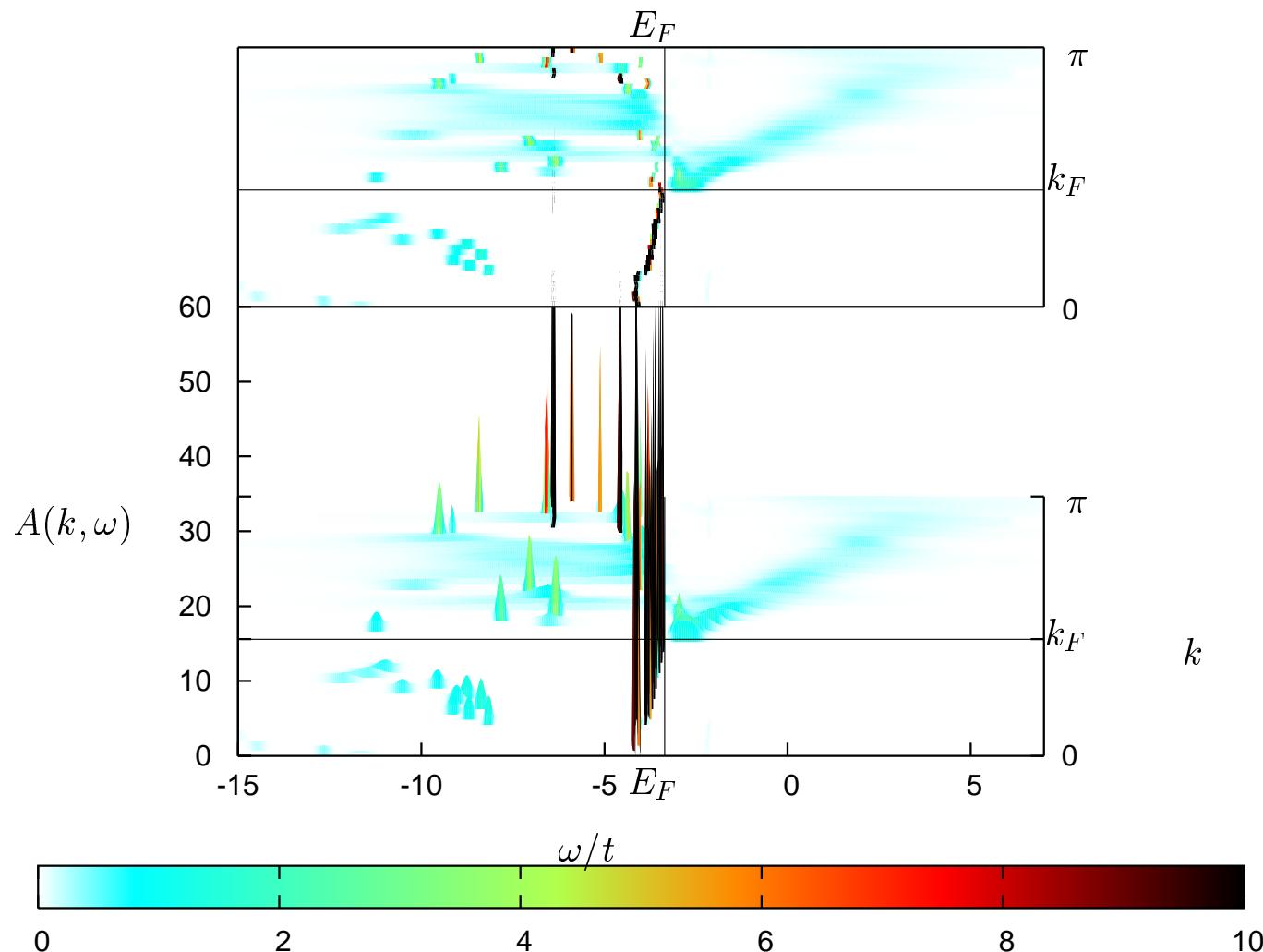
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- Can they Bose-condense for some J and n?
- Can they help to understand $D = 2$?

Photoemission and inverse photoemission with phase separation

$$n = 0.9 \ J/t = 4$$



Discontinuous spectrum on the photoemission side

5.4 Outlook of on-going and future work

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- Spin and density correlation and spectral functions.
→ spin-gap and phase separation

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- Feasibility of two-dimensional simulations
- Extension to finite temperature

New representation and the minus-sign problem

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Minus sign problem in one dimension: non-existent for $N_{\uparrow} + N_{\downarrow} = 4m + 2$,
with m integer

New representation and the minus-sign problem

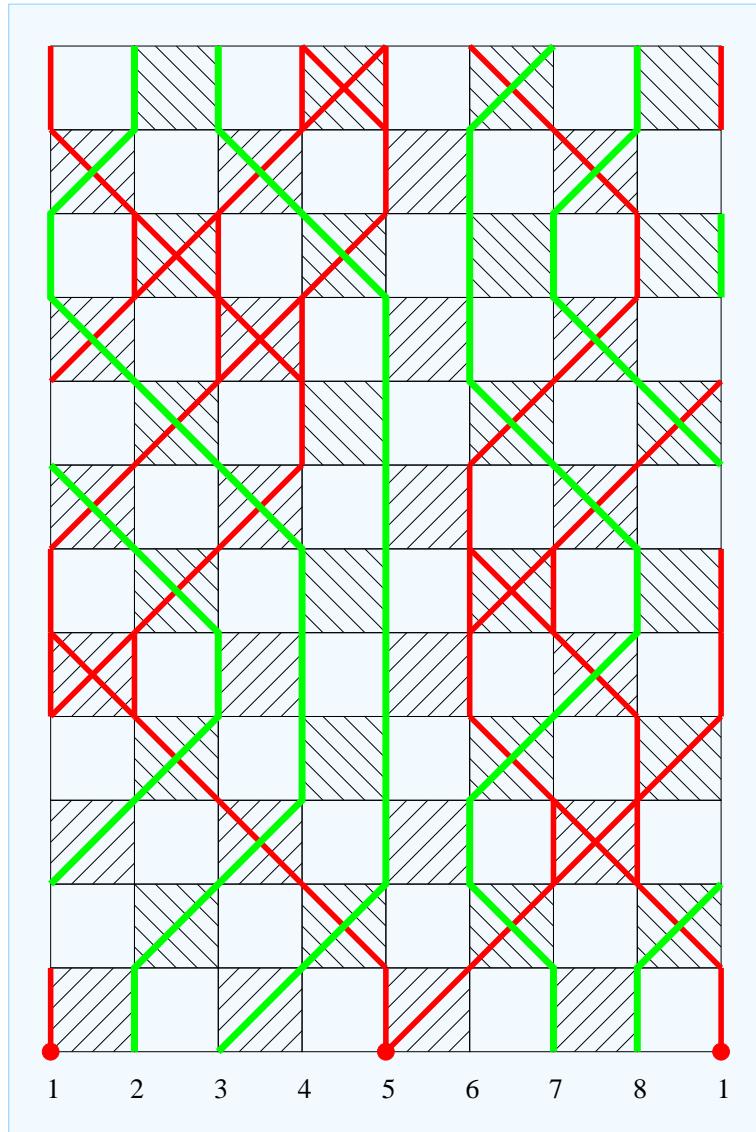
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Observation: $\langle \text{sign} \rangle > 0.95$ for $N_{\uparrow} + N_{\downarrow} = 4m$

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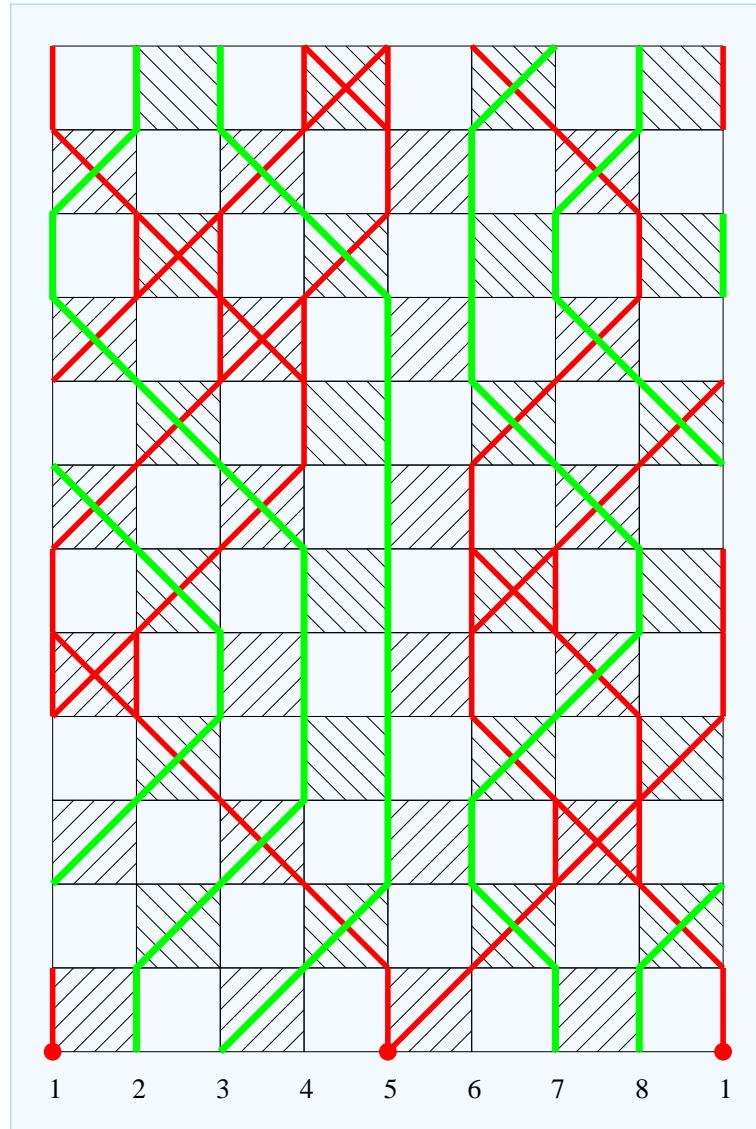
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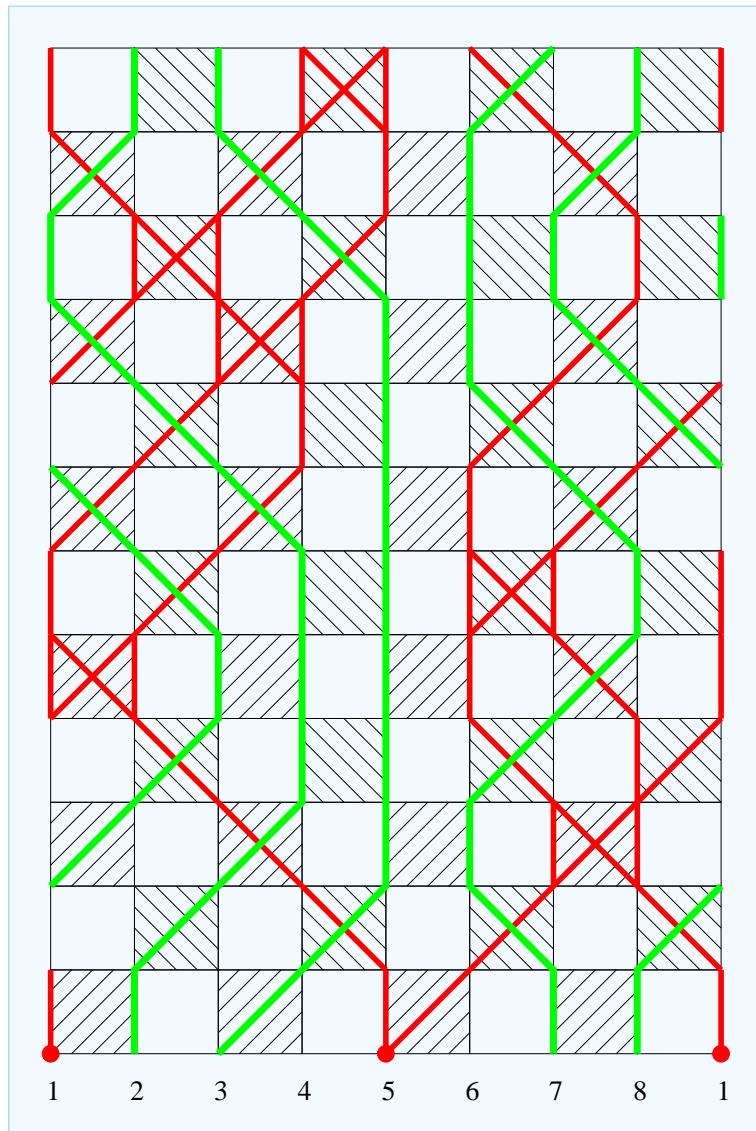


|↓> is a barrier for holes

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$| \downarrow \rangle$ is a barrier for holes

expect also less severe minus sign problem for J small, low doping, and large systems in two dimensions

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- **Checkerboard decomposition** → reduction to a 2-sites problem

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- **Fast Fourier transformation** → alternative to checkerboard decomposition.
- **Jordan-Wigner transformation** → from bosons to fermions in 1-D
For 2-D see

E. Fradkin, Phys. Rev. Lett. 63, 322 (1989)

Y.R. Wang, Phys. Rev. B 43, 3786 (1991)

Summary of Lecture 2

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- Loop-algorithm → particular form of cluster algorithms.
Statistical mechanics of graphs

C. Fortuin and P. Kasteleyn, Physica **57**, 536 (1972)

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S-1/2 Heisenberg antiferromagnet up to $10^3 \times 10^3$ sites

J.-K. Kim and M. Troyer, Phys. Rev. Lett. **80**, 2705 (1998)

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J.-K. Kim and M. Troyer, Phys. Rev. Lett. **80**, 2705 (1998)
- Minus-sign problem → eliminated in special cases with the Meron method

Summary of Lecture 3

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- t-J model → strongly correlated fermions
- Canonical transformation → pseudospins + spinless holes
- Single hole dynamics and loop-algorithm → exact treatment of single hole dynamics
- Single hole in a 2-D quantum antiferromagnet
 - coherent quasiparticle with internal dynamics
 - holons and spinons confined by string potential
 - Self-consistent Born approximation agrees very well with QMC

Summary of Lecture 4

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- Hybrid-loop algorithm → new algorithm for doped antiferromagnets

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- Static and dynamical correlation functions → spectral functions

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- t-J model in 1-D → rich phase diagram reminiscent of high T_c superconductors
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- One-particle spectral functions → QMC for the nearest neighbor t-J model vs. analytical results for the $1/r^2$ t-J model
- Antiholon → new generic excitation of the nearest neighbor t-J model

Collaborators

- Dr. Michael Brunner (HypoVereinsbank-München)
- Dr. Catia Lavalle (University of Stuttgart)
- Dr. Sylvain Capponi (Université Paul Sabatier, Toulouse)
- Dr. Mitsuhiro Arikawa (University of Stuttgart)
- Prof. Dr. Fakher F. Assaad (University of Würzburg)