# Monte Carlo simulations of quantum systems with global updates

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# Quantum spin-systems II

### **Loop-algorithm and further developements**



H.G. Evertz, Adv. Phys. 52, 1 (2003)

H.G. Evertz, Adv. Phys. **52**, 1 (2003)

# Weight of a configuration $\boldsymbol{s} = (s_1, \ldots, s_{2L})$

W(s) (1)

H.G. Evertz, Adv. Phys. **52**, 1 (2003)

Weight of a configuration 
$$s = (s_1, \dots, s_{2L})$$
  

$$W(s) = \sum_{\mathcal{G}} V(\mathcal{G}) \Delta(s, \mathcal{G})$$
(1)

 $V(\mathcal{G}) \longrightarrow$  weight of graph  $\mathcal{G}$ .

$$\Delta(\boldsymbol{s}, \mathcal{G}) = \begin{cases} 1 & \text{if graph } \mathcal{G} \text{ compatible with } \boldsymbol{s} \\ 0 & \text{otherwise }. \end{cases}$$

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Assume (1) is also fulfilled at each plaquette.

$$\hookrightarrow w(u) = \sum_{g} v(g) \Delta(u,g) \;,$$

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Probability of a graph given a configuration on a plaquette

$$p(g \mid u) = \frac{v(g)\Delta(u,g)}{w(u)},$$

#### **Consider all possible configurations of shaded plaquettes**



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#### **Possible graphs**







$$w(u) = \sum_{g} v(g) \Delta(u,g)$$



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$$\downarrow$$
$$e^{-\Delta\tau J\Delta/4}$$
$$= v(\parallel) + v(\times) + v_1(\otimes)$$

$$\sinh(\Delta \tau J/2) e^{\Delta \tau J \Delta/4}$$
$$= v(=) + v(\times) + v_2(\otimes)$$

$$\cosh(\Delta \tau J/2) e^{\Delta \tau J \Delta/4}$$
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 $v(\parallel) = \frac{1}{2} \exp\left(-\frac{\Delta \tau J \Delta}{4}\right) \left\{1 + \exp\left[-\Delta \tau \left(\frac{J}{2} - \frac{J \Delta}{2}\right)\right]\right\},$   
 $v(=) = \frac{1}{2} \exp\left(-\frac{\Delta \tau J \Delta}{4}\right) \left\{\exp\left[\Delta \tau \left(\frac{J}{2} + \frac{J \Delta}{2}\right)\right] - 1\right\}.$ 

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 $\hookrightarrow$  need only two graphs



































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#### **Detailed balance**

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Transition probability (heat-bath)

$$p(\boldsymbol{s} \rightarrow \boldsymbol{s}', \mathcal{G}) = rac{W(\boldsymbol{s}', \mathcal{G})}{W(\boldsymbol{s}, \mathcal{G}) + W(\boldsymbol{s}', \mathcal{G})},$$
### Update

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## All $S_T^z$ states accessible

Simulation in grand canonical ensemble



# → Simulation in grand canonical ensemble







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World lines

**Active loops** 





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**Estimator for the expectation value** 

$$<\mathcal{O}>~=~rac{\sum_{m{s}}W\left(m{s}
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**Estimator for the expectation value** 

$$<\mathcal{O}> = \frac{\sum_{\boldsymbol{s}} W\left(\boldsymbol{s}\right) \ \mathcal{O}\left(\boldsymbol{s}\right)}{\sum_{\boldsymbol{s}} W\left(\boldsymbol{s}\right)}$$

Recall

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Then, we can define

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 $< \mathcal{O} >$  is an expectation value in the ensemble of graphs.

$$4S_i^z S_j^z \left( \mathcal{G} = \text{loop} \right) = \begin{cases} \sigma_i \sigma_j & \text{in the same cluster} \\ 0 & \text{otherwise} \end{cases}$$

where  $\sigma_i = \pm 1$ .

$$4S_i^z S_j^z \left( \mathcal{G} = \text{loop} \right) = \begin{cases} \sigma_i \sigma_j & \text{in the same cluster} \\ 0 & \text{otherwise} \end{cases}$$

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Only pairs i and j in the same loop contribute.

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Recall



vertical lines join sites with equal spins

horizontal lines join sites with alternating spins

$$\hookrightarrow (-1)^{|x_i - x_j|} 4S_i^z S_j^z$$

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Improved estimators reduce fluctuations

## 2.3 Off-diagonal operators

R. Brower, S. Chandrasekharan, and U.J. Wiese, Physica A 261, 520 (1998)
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Recall

$$\langle \mathcal{O} \rangle = \sum_{\mathcal{G}} P(\mathcal{G}) \mathcal{O}(\mathcal{G})$$

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 $\hookrightarrow \mathcal{O}\left(\mathcal{G}\right) \text{ is basis independent}$ 

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 $\hookrightarrow$  Change to a basis where  ${\mathcal O}$  is diagonal

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 $\hookrightarrow \mathcal{O}\left(\mathcal{G}\right) \text{ is basis independent}$ 

 $\hookrightarrow$  Change to a basis where  ${\mathcal O}$  is diagonal

As in the case of diagonal operators,

$$4S_i^x S_j^x \left( \mathcal{G} = \text{loop} \right) = \begin{cases} \sigma_i \sigma_j & \text{in the same cluster} \\ 0 & \text{otherwise} \end{cases}$$

B.B. Beard and U.J. Wiese, Phys. Rev. Lett. 77, 5130 (1996)

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**Discrete imaginary time**  $\longrightarrow$  **systematic error**  $\mathcal{O}\left(\Delta\tau^2\right)$ 

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Take the limit  $\Delta \tau \rightarrow 0$ 

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#### Take the limit $\Delta \tau \rightarrow 0$





 $p(\|) = 1$ 







## **Continuum limit** $\longrightarrow$ probability density per unit time

$$\frac{p(=)}{\Delta \tau} \longrightarrow \frac{J}{2}$$



# Start at some point of the lattice



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Start at some point of the lattice

Put an horizontal line with probability density  $\rho(=)=J/2$ 



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**Reverse direction** 



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**Repeat until loop closes** 

Flip the loop







#### • Infinite lattices and zero temperature

H.G. Evertz and W. von der Linden, Phys. Rev. Lett. 86, 5164 (2001)

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## • Hubbard model

N. Kawashima, J.E. Gubernatis, and H.G. Evertz, Phys. Rev. B 50, 136 (1994)

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#### • t-J model

B. Ammon, H.G. Evertz, N. Kawashima, M. Troyer, and B. Frischmuth,
Phys. Rev. B 58, 4304 (1998)
M. Brunner and A. Muramatsu, Phys. Rev. B 58, R10100 (1998)

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#### Solution of the minus sign problem with merons

S. Chandrasekharan and U.J. Wiese, Phys. Rev. Lett. 83, 3116 (1999)

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#### Solution of the minus sign problem with merons

S. Chandrasekharan and U.J. Wiese, Phys. Rev. Lett. 83, 3116 (1999)

#### Stochastic series expansion

O.F. Syljuasen and A.W. Sandvik, Phys. Rev. E 66, 046701 (2002)

# Loop-algorithm for the t-J model Two dimensional t-model (J=0) with one hole

## Nagaoka's ferromagnetism



M. Brunner and A. Muramatsu, Phys. Rev. B 58, R10100 (1998)